



# Breaking the degeneracy between adiabatic and impulsive cusp-core transformation mechanisms

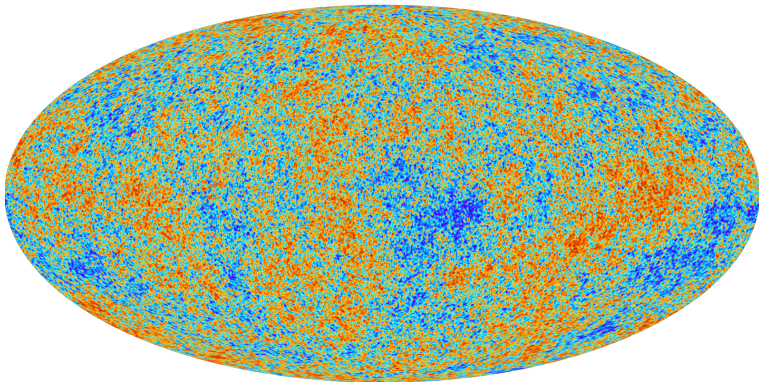
Jan David Burger

Faculty of Physical Sciences  
University of Iceland

2021

# Structure formation in $\Lambda$ CDM

■ CMB:  $\frac{T - \bar{T}}{\bar{T}} \sim 10^{-5}$



ESA/Planck

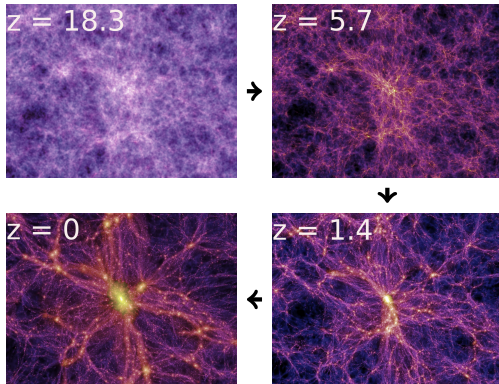


# Structure formation in $\Lambda$ CDM

- CMB:  $\frac{T-\bar{T}}{\bar{T}} \sim 10^{-5}$
- Today: local density contrast  
 $\frac{\rho-\bar{\rho}}{\bar{\rho}} \gg 1$

# Structure formation in $\Lambda$ CDM

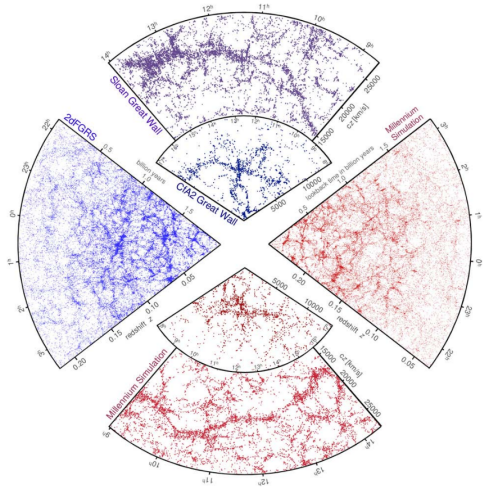
- CMB:  $\frac{T-\bar{T}}{\bar{T}} \sim 10^{-5}$
- Today: local density contrast  $\frac{\rho-\bar{\rho}}{\bar{\rho}} \gg 1$
- Growth of perturbations linear until  $\frac{\rho-\bar{\rho}}{\bar{\rho}} \lesssim 1$
- $N$ -body simulations follow gravitational collapse in the non-linear regime



Springel et. al. 2005

# Structure formation in $\Lambda$ CDM

- CMB:  $\frac{T-\bar{T}}{\bar{T}} \sim 10^{-5}$
- Today: local density contrast  $\frac{\rho-\bar{\rho}}{\bar{\rho}} \gg 1$
- Growth of perturbations linear until  $\frac{\rho-\bar{\rho}}{\bar{\rho}} \lesssim 1$
- $N$ -body simulations follow gravitational collapse in the non-linear regime
- Simulated structure matches observations on large scales



Springel et. al. 2006

# Challenges on small scales

<https://wwwmpa.mpa-garching.mpg.de/aquarius/>

$$r_{200} \sim 200 \text{ kpc}$$

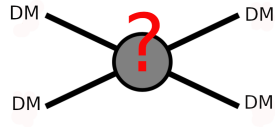
$$d \sim 60 \text{ kpc}$$

<https://solarsystem.nasa.gov/resources/285/the-milky-way-galaxy/>

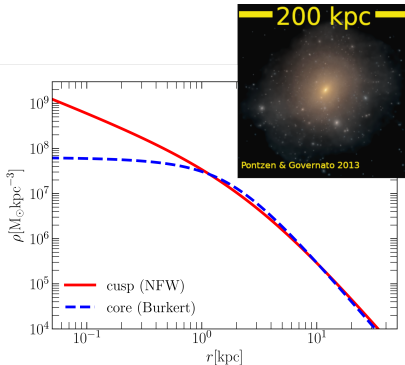
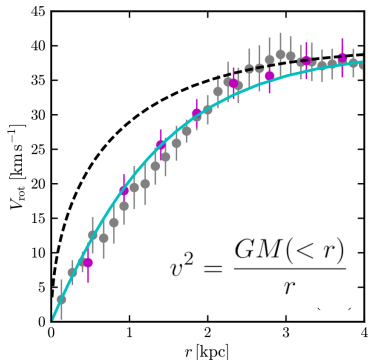
## Baryonic physics



## DM particle physics



# Challenges on small scales



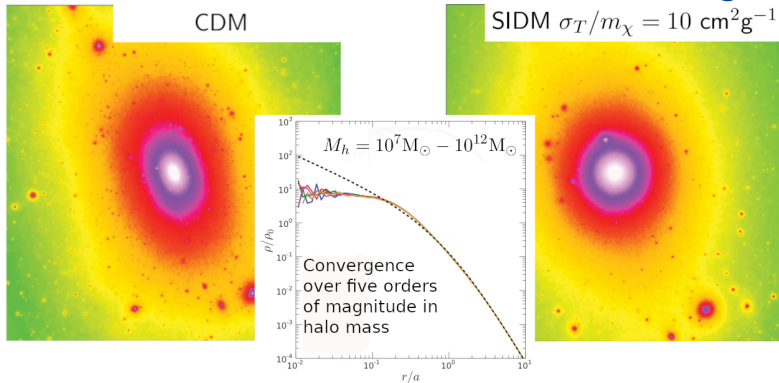
Bullock & Boylan-Kolchin 2017

- Rotation curves of some dwarf galaxies rise slower than expected
- Their host DM haloes appear to be cored instead of cuspy

# Scope of this presentation

- Inner density profile of simulated CDM haloes has a universal, cuspy shape
- Measured rotation curves of dwarf galaxies suggest DM haloes with constant density cores
- Theory and observation need to be reconciled by core formation mechanisms
- Different core formation mechanisms must lead to (nearly) identical, *i.e. degenerate* final halo density profiles
- The timescale of the cusp-core transformation is different for different feasible mechanisms
- Key motivation: Can we use this difference to break the degeneracy between core formation mechanisms?

# Core formation mechanisms - Self-interacting DM



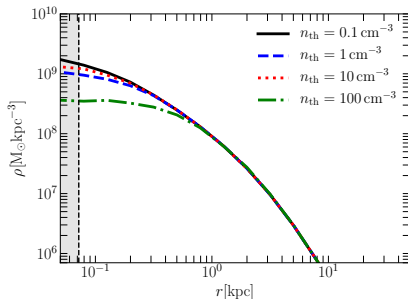
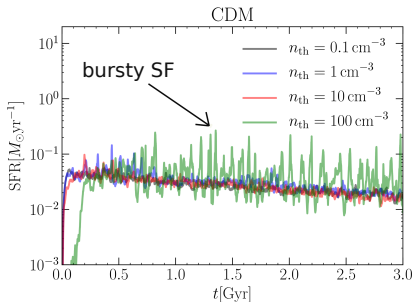
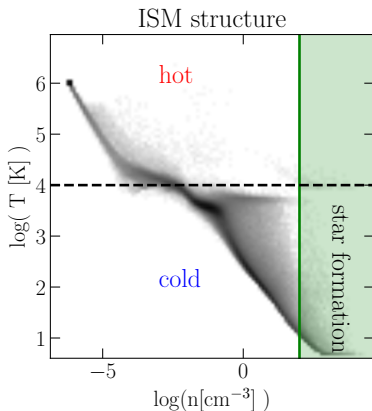
Vogelsberger et al. 2012

- Elastic scattering between DM particles  $\rightarrow$  heat exchange
- Scattering rate regulated by momentum transfer cross section:

$$P_{ij} \propto \rho_i \frac{\sigma_T}{m_\chi} v_{ij} \Delta t_i,$$

# Core formation mechanisms - SN Feedback

- Bursty star formation causes impulsive SNF
- Core formation if central potential is dominated by baryons





# Analytic models for SNF

- We can model the impact of mass loss without running full hydrodynamical simulations

# Analytic models for SNF

- We can model the impact of mass loss without running full hydrodynamical simulations
- One-component effective model for SNF in DM only simulations:
  - ▶ Periodically add and remove central potential
  - ▶ Potential and injection scheme define model

# Analytic models for SNF

- We can model the impact of mass loss without running full hydrodynamical simulations
- One-component effective model for SNF in DM only simulations:
  - ▶ Periodically add and remove central potential
  - ▶ Potential and injection scheme define model
- Two-component effective model for SNF in DM only simulations:
  - ▶ External potential mimicking a galaxy
  - ▶ Scheme to model the formation of “superbubbles”
  - ▶ Distribution of “superbubbles” determined by external potential

# Analytic models for SNF

- We can model the impact of mass loss without running full hydrodynamical simulations
- One-component effective model for SNF in DM only simulations:
  - ▶ Periodically add and remove central potential
  - ▶ Potential and injection scheme define model
- Two-component effective model for SNF in DM only simulations:
  - ▶ External potential mimicking a galaxy
  - ▶ Scheme to model the formation of “superbubbles”
  - ▶ Distribution of “superbubbles” determined by external potential
- Free parameters of two-component model:
  - ▶ Galaxy type and galaxy size
  - ▶ Energy injection time
  - ▶ Total amount of injected energy

# Adiabatic and impulsive evolution

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

where the actions are conserved quantities.

# Adiabatic and impulsive evolution

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

where the actions are conserved quantities.

- In time-dependent spherical potentials, the evolution of a **tracer particle's** radial action depends on how fast the potential changes:

# Adiabatic and impulsive evolution

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

where the actions are conserved quantities.

- In time-dependent spherical potentials, the evolution of a **tracer particle's** radial action depends on how fast the potential changes:
  - ▶ In **adiabatically** evolving potentials (evolution slow compared with typical orbital times), actions are conserved

# Adiabatic and impulsive evolution

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

where the actions are conserved quantities.

- In time-dependent spherical potentials, the evolution of a **tracer particle's** radial action depends on how fast the potential changes:
  - ▶ In **adiabatically** evolving potentials (evolution slow compared with typical orbital times), actions are conserved
  - ▶ In **impulsively** evolving potentials, actions of tracers change discontinuously (and their energies change by an explicitly phase-dependent amount)



# Adiabatic and impulsive evolution

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

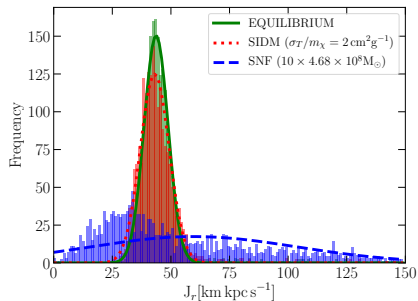
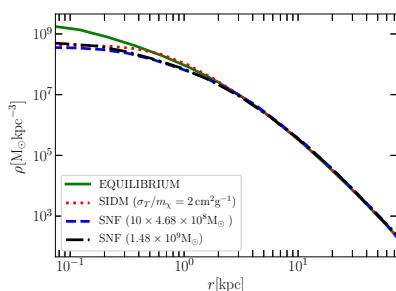
where the actions are conserved quantities.

- In time-dependent spherical potentials, the evolution of a **tracer particle's** radial action depends on how fast the potential changes:
  - ▶ In **adiabatically** evolving potentials (evolution slow compared with typical orbital times), actions are conserved
  - ▶ In **impulsively** evolving potentials, actions of tracers change discontinuously (and their energies change by an explicitly phase-dependent amount)

## Key idea

Adiabatic /impulsive core formation mechanisms may affect the kinematic properties of stars in distinct ways.

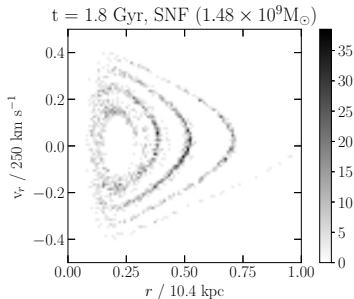
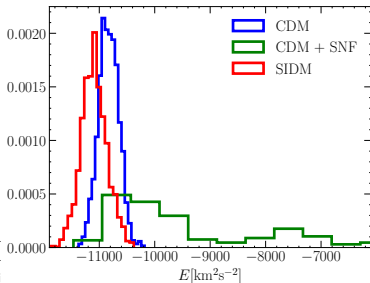
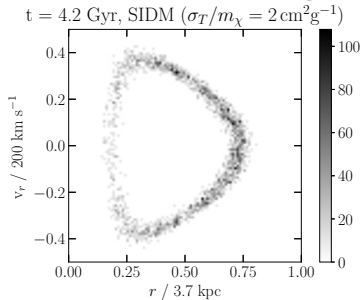
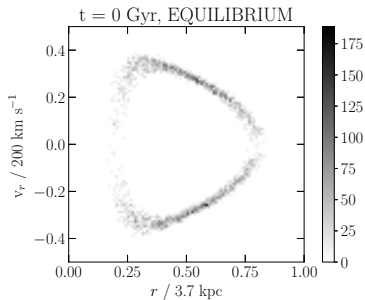
# Core formation through SIDM and SNF



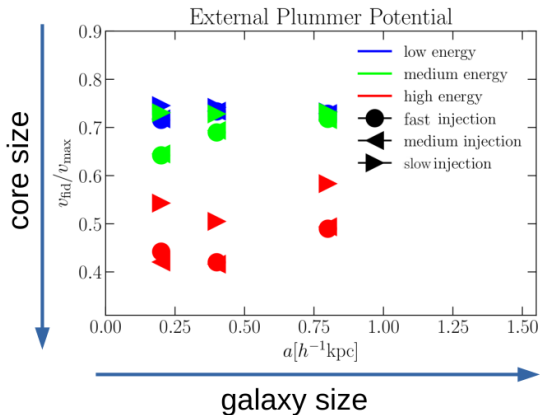
Burger & Zavala 2019

- $10^7$  DM particles,  $M_{200} \sim 1.5 \times 10^{10} M_{\odot}$ ,  $c_{200} = 15$
- One-component model for SNF, Probabilistic model for SIDM

# Orbital families and phase mixing



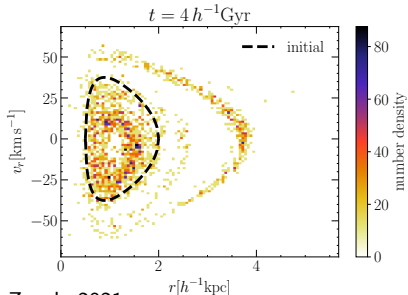
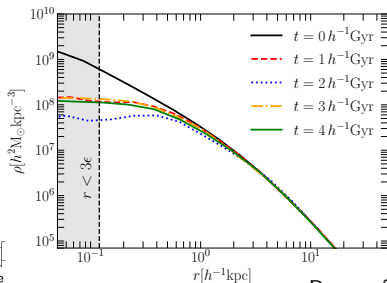
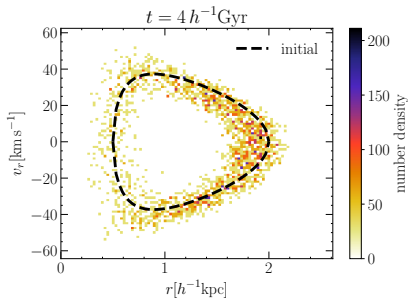
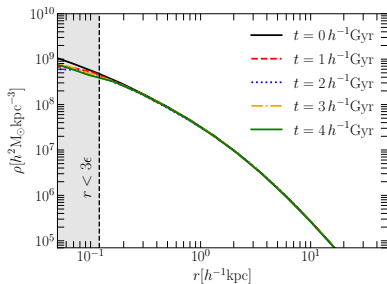
# Evolution of circular velocity curves



Burger & Zavala 2021

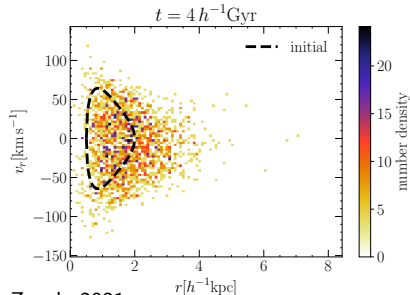
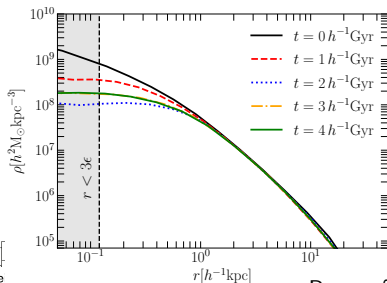
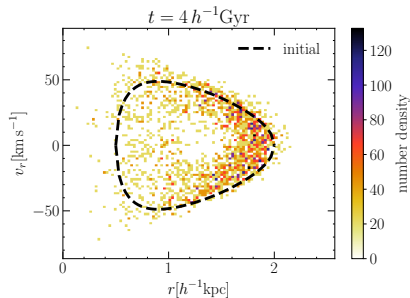
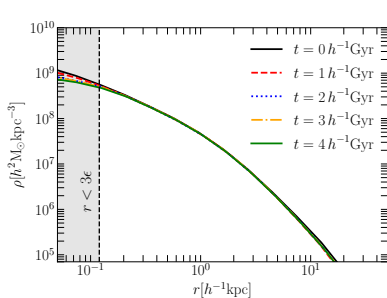
- Two-component model, initially  $M_h = 10^{10} h^{-1} M_{\odot}$   $c_{200} = 13$
- Galaxy size, injected energy and injection time affect core size
- Similar picture for Plummer sphere and disk

# Orbital families and symmetry: Plummer sphere



Burger & Zavala 2021

# Orbital families and symmetry: Flat disk



Burger & Zavala 2021

# Hydrodynamic simulations of an isolated dwarf

## ■ Set up ICs of an isolated SMC-analog in a live DM halo

- ▶  $M_h \sim 1.6 \times 10^{10} M_\odot$        $c_{200} \sim 17$
- ▶ Gas disk, stellar disk & stellar bulge  $\sim 5\%$  of halo mass
- ▶  $R_\star = 0.7 \text{ kpc}$        $R_{\text{gas}} = 2.1 \text{ kpc}$
- ▶ Gas with  $Z = Z_\odot$ ,  $T = 10^4 \text{ K}$ , and in hydrostatic equilibrium
- ▶ Particle / gas cell mass  $\sim 10^3 M_\odot$

# Hydrodynamic simulations of an isolated dwarf

- Set up ICs of an isolated SMC-analog in a live DM halo
  - ▶  $M_h \sim 1.6 \times 10^{10} M_\odot$       $c_{200} \sim 17$
  - ▶ Gas disk, stellar disk & stellar bulge  $\sim 5\%$  of halo mass
  - ▶  $R_\star = 0.7 \text{ kpc}$       $R_{\text{gas}} = 2.1 \text{ kpc}$
  - ▶ Gas with  $Z = Z_\odot$ ,  $T = 10^4 \text{ K}$ , and in hydrostatic equilibrium
  - ▶ Particle / gas cell mass  $\sim 10^3 M_\odot$
- Isolated galaxy evolved using SMUGGLE, which models
  - ▶ Cooling and heating of gas
  - ▶ Star formation and stellar evolution and feedback



# Hydrodynamic simulations of an isolated dwarf

- Set up ICs of an isolated SMC-analog in a live DM halo
  - ▶  $M_h \sim 1.6 \times 10^{10} M_\odot$        $c_{200} \sim 17$
  - ▶ Gas disk, stellar disk & stellar bulge  $\sim 5\%$  of halo mass
  - ▶  $R_\star = 0.7 \text{ kpc}$        $R_{\text{gas}} = 2.1 \text{ kpc}$
  - ▶ Gas with  $Z = Z_\odot$ ,  $T = 10^4 \text{ K}$ , and in hydrostatic equilibrium
  - ▶ Particle / gas cell mass  $\sim 10^3 M_\odot$
- Isolated galaxy evolved using SMUGGLE, which models
  - ▶ Cooling and heating of gas
  - ▶ Star formation and stellar evolution and feedback
- Simulations include elastic scatters between DM particles

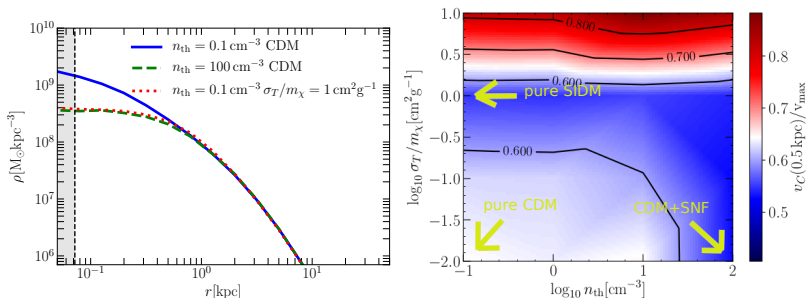
# Hydrodynamic simulations of an isolated dwarf

- Set up ICs of an isolated SMC-analog in a live DM halo
  - ▶  $M_h \sim 1.6 \times 10^{10} M_\odot$        $c_{200} \sim 17$
  - ▶ Gas disk, stellar disk & stellar bulge  $\sim 5\%$  of halo mass
  - ▶  $R_\star = 0.7 \text{ kpc}$        $R_{\text{gas}} = 2.1 \text{ kpc}$
  - ▶ Gas with  $Z = Z_\odot$ ,  $T = 10^4 \text{ K}$ , and in hydrostatic equilibrium
  - ▶ Particle / gas cell mass  $\sim 10^3 M_\odot$
- Isolated galaxy evolved using SMUGGLE, which models
  - ▶ Cooling and heating of gas
  - ▶ Star formation and stellar evolution and feedback
- Simulations include elastic scatters between DM particles
- Systematically vary  $n_{\text{th}}$  and  $\sigma_T/m_\chi$  in suite of 16 simulations

# Hydrodynamic simulations of an isolated dwarf

- Set up ICs of an isolated SMC-analog in a live DM halo
  - ▶  $M_h \sim 1.6 \times 10^{10} M_\odot$      $c_{200} \sim 17$
  - ▶ Gas disk, stellar disk & stellar bulge  $\sim 5\%$  of halo mass
  - ▶  $R_\star = 0.7 \text{ kpc}$      $R_{\text{gas}} = 2.1 \text{ kpc}$
  - ▶ Gas with  $Z = Z_\odot$ ,  $T = 10^4 \text{ K}$ , and in hydrostatic equilibrium
  - ▶ Particle / gas cell mass  $\sim 10^3 M_\odot$
- Isolated galaxy evolved using SMUGGLE, which models
  - ▶ Cooling and heating of gas
  - ▶ Star formation and stellar evolution and feedback
- Simulations include elastic scatters between DM particles
- Systematically vary  $n_{\text{th}}$  and  $\sigma_T/m_\chi$  in suite of 16 simulations
- Goal: break degeneracy between adiabatically and impulsively formed cores with kinematic properties of stars / gas

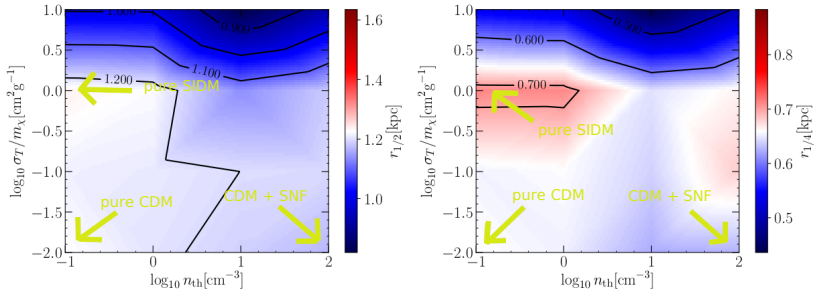
# Degeneracy of DM density profiles



Burger et al. 2021

Different combinations of  $n_{\text{th}}$  and  $\sigma_T/m_\chi$  can produce cores of a similar size!

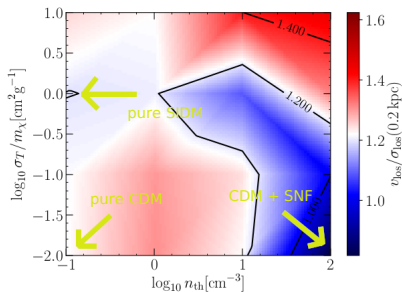
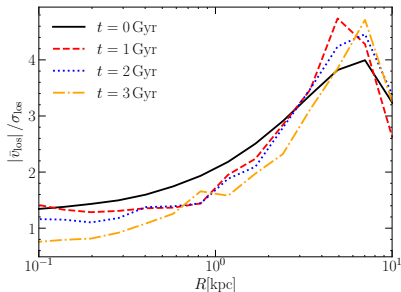
# Breaking the degeneracy



Burger et al. 2021

- Adiabatic core formation causes migration of stellar orbits

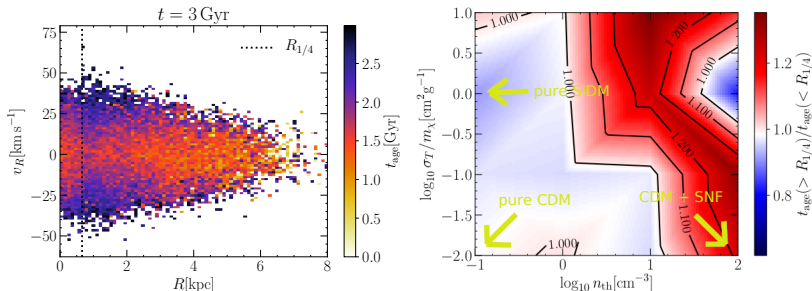
# Breaking the degeneracy



Burger et al. 2021

- Adiabatic core formation causes migration of stellar orbits
- Impulsive SNF can lead to an increase in random gas motion

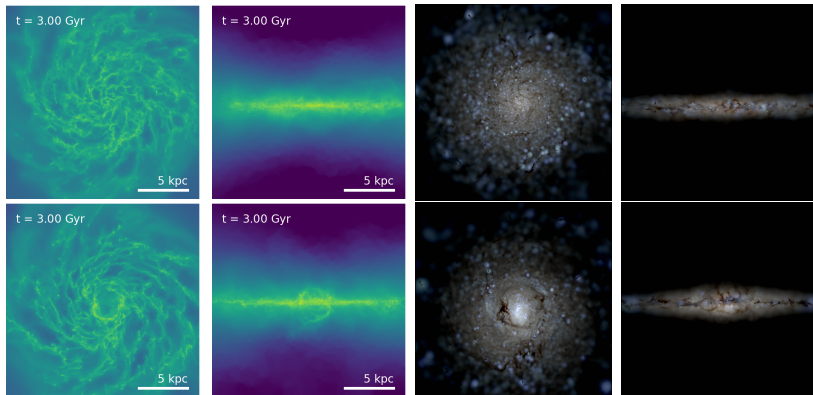
# Breaking the degeneracy



Burger et al. 2021

- Adiabatic core formation causes migration of stellar orbits
- Impulsive SNF can lead to an increase in random gas motion
- Impulsive SNF can cause large stellar age gradients

# CDM runs – low vs high star formation threshold

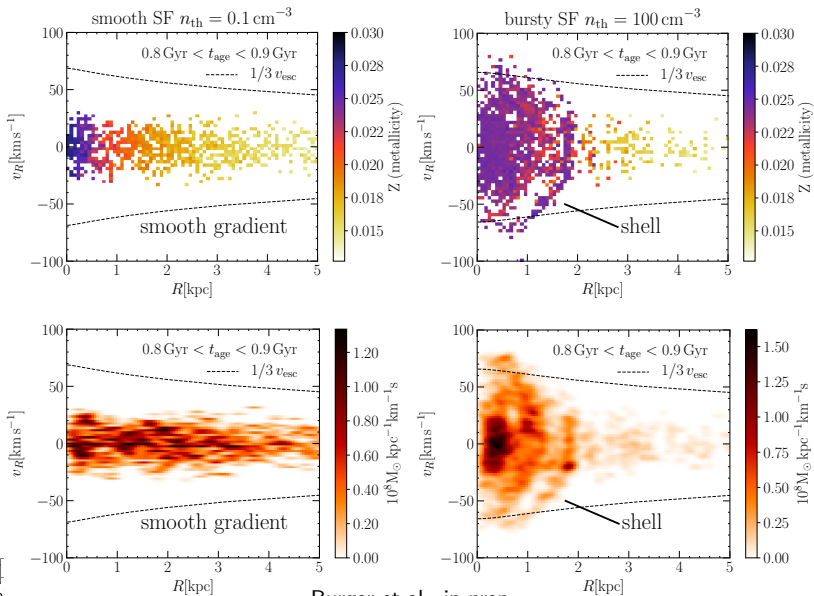


Burger et al., in prep.

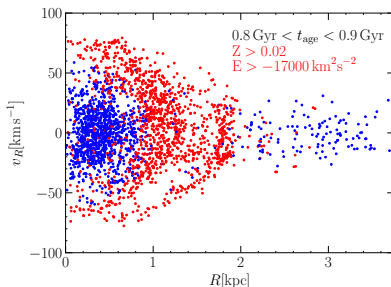
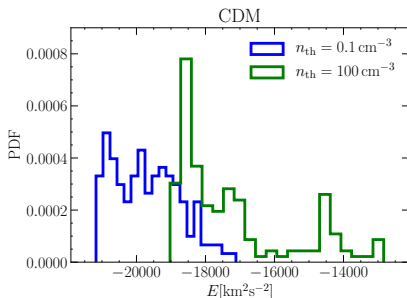
- Compare CDM runs with  $n_{\text{th}} = 0.1 \text{ cm}^{-3}$  and  $n_{\text{th}} = 100 \text{ cm}^{-3}$
- Supperbubble created by impulsive SNF is apparent in the edge-on projection of the gas



# Phase space distribution of mono-age stars



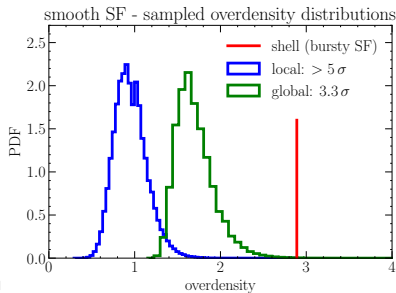
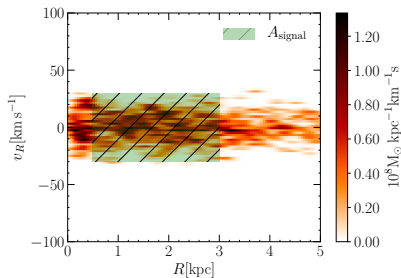
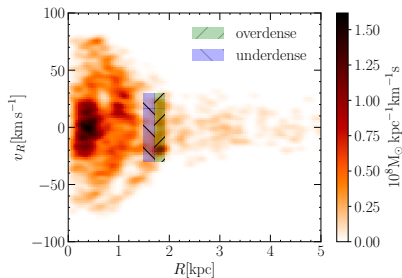
# Explaining the shell feature



Burger et al., in prep.

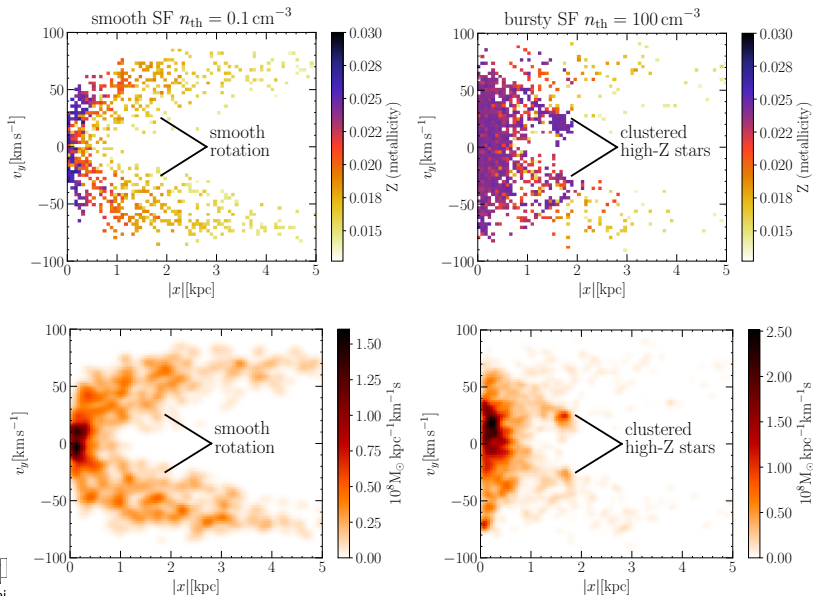
- Stars of similar age and metallicity form orbital families
- Look at 150 most metal-rich stars
- Impulsive SNF causes discontinuous energy increase
- High-energy particles make up the shell → phase mixing

# Significance of the shell feature



- Use smooth distribution as target distribution for random sampling
- Determine likelihood of overdensity being a chance occurrence
- Feature has local (global) significance of  $> 5\sigma$  ( $3.3\sigma$ )

# LOS projection of phase space shell



# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies

# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores

# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores
- This degeneracy is broken by the kinematic properties of stars and gas

# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores
- This degeneracy is broken by the kinematic properties of stars and gas
- SIDM-induced core formation causes an adiabatic expansion of central stars, while impulsive SNF creates large age gradients



# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores
- This degeneracy is broken by the kinematic properties of stars and gas
- SIDM-induced core formation causes an adiabatic expansion of central stars, while impulsive SNF creates large age gradients
- Shell-like features associated with early-stage phase mixing appear in the phase space distribution of mono-age stars in the aftermath of impulsive SNF events

# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores
- This degeneracy is broken by the kinematic properties of stars and gas
- SIDM-induced core formation causes an adiabatic expansion of central stars, while impulsive SNF creates large age gradients
- Shell-like features associated with early-stage phase mixing appear in the phase space distribution of mono-age stars in the aftermath of impulsive SNF events
- We need cosmological simulations to further quantify the signal(s) and take environmental effects into account

# Conclusion and outlook

- Both SIDM and SNF are viable cusp-core transformation mechanisms at the scale of dwarf galaxies
- They can create nearly indistinguishable DM cores
- This degeneracy is broken by the kinematic properties of stars and gas
- SIDM-induced core formation causes an adiabatic expansion of central stars, while impulsive SNF creates large age gradients
- Shell-like features associated with early-stage phase mixing appear in the phase space distribution of mono-age stars in the aftermath of impulsive SNF events
- We need cosmological simulations to further quantify the signal(s) and take environmental effects into account
- Required observational data is likely available 5-10 years down the road (e.g. Roman Space Telescope)

# References

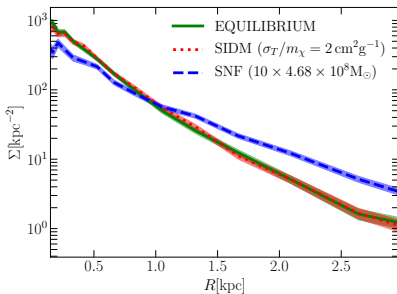
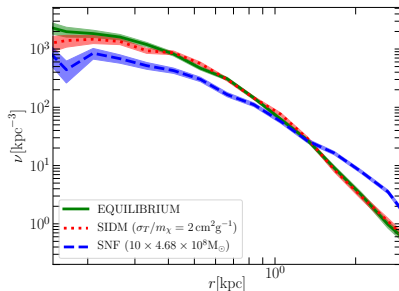
This presentation contains results from four research articles:

- *The nature of core formation in dark matter haloes: adiabatic or impulsive?* by Jan D. Burger and Jesús Zavala, 1810.10024
- *SN-driven mechanism of cusp-core transformation: an appraisal* by Jan D. Burger and Jesús Zavala, 2103.01231
- *Degeneracies Between Self-interacting Dark Matter and Supernova Feedback as cusp-core transformation mechanisms* by Jan D. Burger, Jesús Zavala, Laura V. Sales, Mark Vogelsberger, Federico Marinacci, and Paul Torrey, 2108.07358
- *Kinematic signatures of impulsive supernova feedback in dwarf galaxies* by Jan D. Burger, Jesús Zavala, Laura V. Sales, Mark Vogelsberger, Federico Marinacci, and Paul Torrey, in prep.

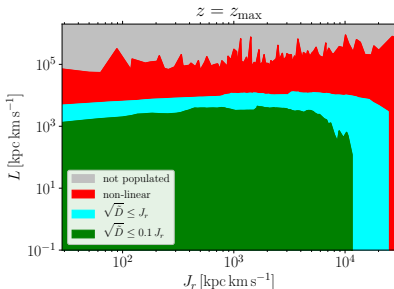
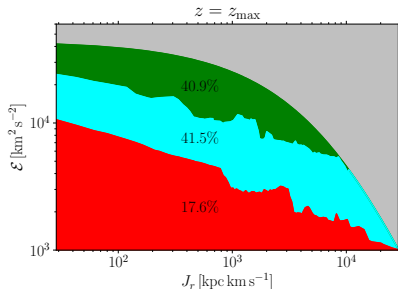
Thank you for your attention !



# Plummer profiles

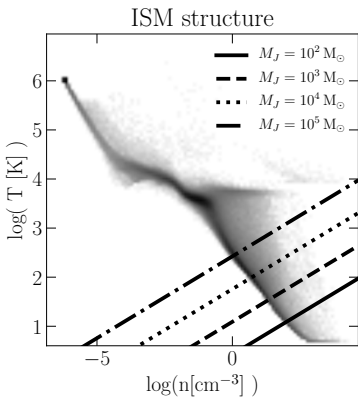
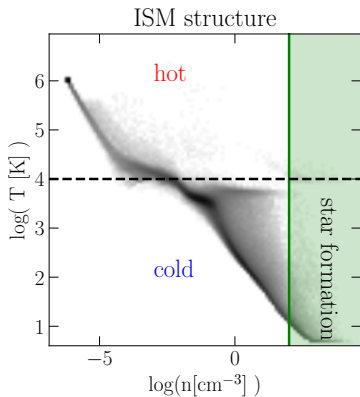


# Action condition in $J_r$ - $L$ space

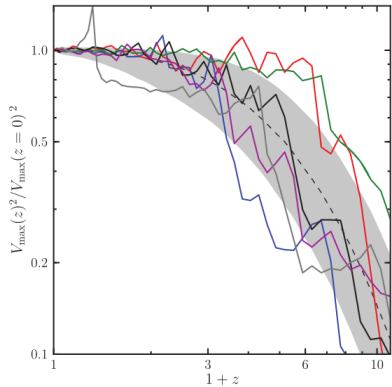
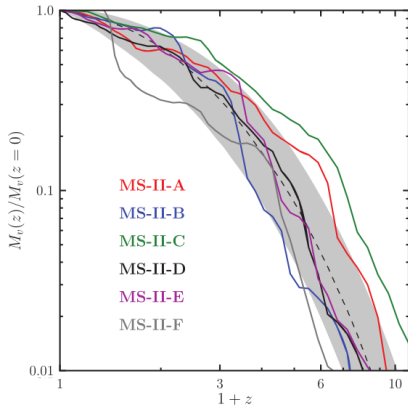




# $n - T$ diagram and Jeans mass

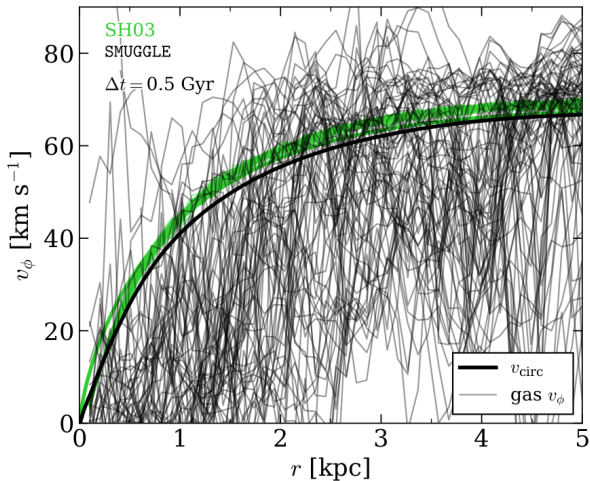


# BK mass accretion history



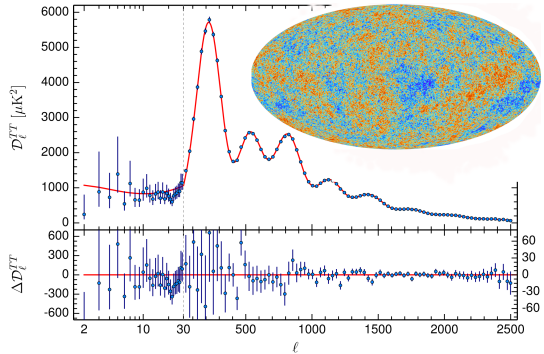
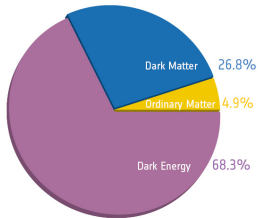
Boylan-Kolchin et al., 2010

# SMUGGLE rotation curves



Jahn et al., in prep.

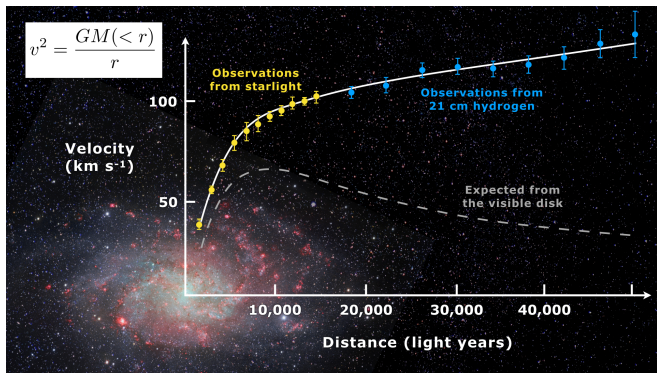
# Evidence for dark matter



ESA/Planck

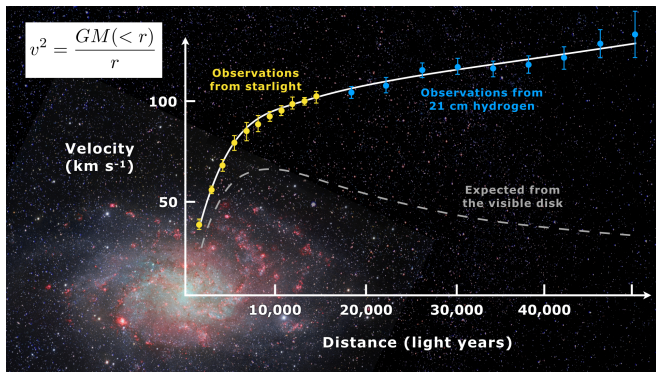
■ Angular power spectrum of the CMB temperature fluctuations

# Evidence for dark matter



- Angular power spectrum of the CMB temperature fluctuations
- The inner kinematics of observed galaxies and galaxy clusters

# Evidence for dark matter



- Angular power spectrum of the CMB temperature fluctuations
- The inner kinematics of observed galaxies and galaxy clusters
- There is also strong evidence for DM from cosmic structure formation

# Radial actions in evolving spherical potentials

- In static potentials with spherical symmetry

$$H(\mathbf{x}, \mathbf{p}) \rightarrow H(J_r, L, L_z),$$

where the actions are conserved quantities

- In mildly time-dependent potentials, the actions of kinematic tracers (nearly massless particles) evolve as

$$J_r = J_{r'} + (\mathbf{r} \cdot \dot{\mathbf{r}}) \frac{\dot{R}}{R} \frac{P(E, L, t)}{2\pi} \equiv J_{r'} + \Delta_{J_r}$$

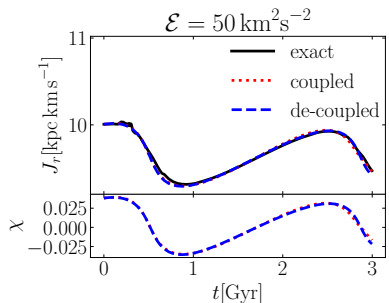
- $\dot{R}/R$  depends on the rate of change in both amplitude and shape of the potential
- $P(E, L, t)/2\pi$  is the reduced radial period

# Radial actions of tracer particles

$$\Phi(r, t) = -\frac{GM_0(1 + \epsilon t)}{r}$$

$$\log_{10}(M_0/M_\odot) = 8$$

$$\epsilon^{-1} = 30 \text{ Gyr}$$

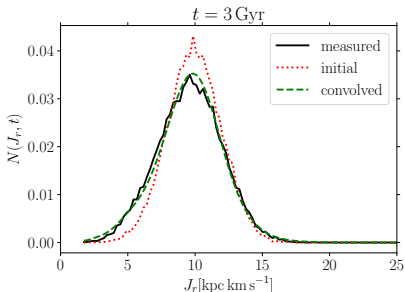
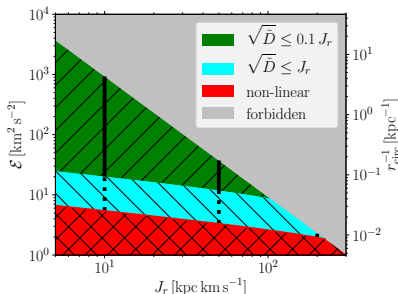


Burger et al. 2020

- The first-order approximation is accurate if the potentials evolves slowly relative to the tracer's orbital time



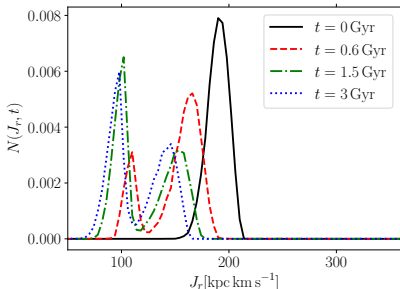
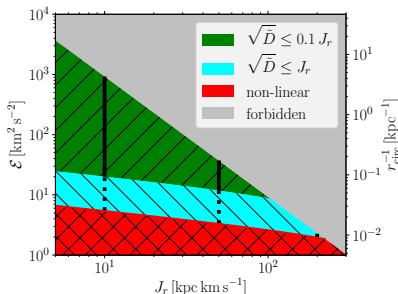
# Radial actions of tracer particles



Burger et al. 2020

- The first-order approximation is accurate if the potentials evolves slowly relative to the tracer's orbital time
- A diffusion theory with drift  $\tilde{C} \sim \langle \Delta_{J_r} \rangle$  and diffusion  $\tilde{D} \sim \langle \Delta_{J_r}^2 \rangle$  accurately predicts the evolution of radial action distributions for  $\tilde{D} \ll J_r^2$

# Radial actions of tracer particles



Burger et al. 2020

- The first-order approximation is accurate if the potentials evolves slowly relative to the tracer's orbital time
- A diffusion theory with drift  $\tilde{C} \sim \langle \Delta_{J_r} \rangle$  and diffusion  $\tilde{D} \sim \langle \Delta_{J_r}^2 \rangle$  accurately predicts the evolution of radial action distributions for  $\tilde{D} \ll J_r^2$ , **but fails for  $\tilde{D} \gtrsim J_r^2$**